

Canonical Ensemble: Calculation of Partition Function



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Calculation of partition function Z of the system

(a) Two level system

The energy of system be either $-\frac{\Delta}{2}$ or $\frac{\Delta}{2}$
then

$$\begin{aligned} Z &= \sum_r \exp(-\beta E_r) \\ &= e^{\beta \frac{\Delta}{2}} + e^{-\beta \frac{\Delta}{2}} \\ &= 2 \cosh\left(\frac{\beta \Delta}{2}\right) \end{aligned}$$

$$\begin{aligned} \text{Internal energy } U &= -\frac{\partial}{\partial \beta} \ln Z \\ &= -\frac{\Delta}{2} \tanh\left(\frac{\beta \Delta}{2}\right) \end{aligned}$$

heat capacity

$$\begin{aligned} C_V &= \left(\frac{\partial U}{\partial T}\right)_V = \left(\frac{\partial U}{\partial \beta}\right)_V \frac{\partial \beta}{\partial T} \\ &= -\left(\frac{\Delta}{2}\right)^2 \operatorname{sech}^2\left(\frac{\beta \Delta}{2}\right) \cdot \left(-\frac{1}{kT^2}\right) \\ &= k \left(\frac{\beta \Delta}{2}\right)^2 \operatorname{sech}^2\left(\frac{\beta \Delta}{2}\right) \end{aligned}$$

Helmholtz free energy

$$F = -kT \ln Z$$

$$= -kT \ln \left\{ 2 \cosh\left(\frac{\beta \Delta}{2}\right) \right\}$$

Entropy

$$S = \frac{U - F}{T}$$

$$= -\frac{\Delta}{2T} \tanh\left(\frac{\beta \Delta}{2}\right) + k \ln \left\{ 2 \cosh\left(\frac{\beta \Delta}{2}\right) \right\}$$

(b) Simple harmonic oscillator

The energy of the system is $(n + \frac{1}{2})\hbar\omega$
where $n = 0, 1, 2, \dots$

$$\begin{aligned} \therefore Z &= \sum_r \exp(-\beta E_r) = \sum_{n=0}^{\infty} \exp\left[-\beta\left(n + \frac{1}{2}\right)\hbar\omega\right] \\ &= e^{-\beta \frac{\hbar\omega}{2}} \cdot \sum_{n=0}^{\infty} \exp(-n\beta\hbar\omega) = \frac{e^{-\beta \frac{\hbar\omega}{2}}}{1 - \exp(-\beta\hbar\omega)} \\ &= \frac{1}{e^{\frac{\beta\hbar\omega}{2}} - e^{-\frac{\beta\hbar\omega}{2}}} \end{aligned}$$

internal energy $U = - \frac{\partial \ln z}{\partial \beta}$

$$\ln z = -\frac{1}{2} \beta \hbar \omega - \ln \{1 - e^{-\beta \hbar \omega}\}$$

$$\therefore U = \frac{1}{2} \hbar \omega + \frac{e^{-\beta \hbar \omega} \cdot \hbar \omega}{1 - e^{-\beta \hbar \omega}}$$

$$= \hbar \omega \left[\frac{1}{2} + \frac{1}{e^{\beta \hbar \omega} - 1} \right]$$

specific heat $C_v = \left(\frac{\partial U}{\partial T} \right)_v$

$$= \hbar \omega \frac{(-1) \hbar \omega}{[e^{\beta \hbar \omega} - 1]^2} e^{\beta \hbar \omega} \left(-\frac{1}{k T^2} \right)$$

$$= k (\beta \hbar \omega)^2 \frac{e^{\beta \hbar \omega}}{[e^{\beta \hbar \omega} - 1]^2}$$

At high temperature $\beta \hbar \omega \ll 1$

$$\therefore e^{\beta \hbar \omega} \approx 1 + \beta \hbar \omega$$

$$\therefore C_v = k \frac{(\beta \hbar \omega)^2 [1 + \beta \hbar \omega]}{(\beta \hbar \omega)^2}$$

$$= k$$

$$\text{and } U \approx \frac{\hbar \omega}{2} + \frac{\hbar \omega}{\beta \hbar \omega} = kT \left[1 + \frac{\hbar \omega}{2kT} \right]$$

$$\approx kT$$

Helmholtz function

$$F = -kT \ln Z$$

$$= -kT \left[-\frac{1}{2} \beta \hbar \omega - \ln \{ 1 - e^{-\beta \hbar \omega} \} \right]$$

$$= \frac{1}{2} \hbar \omega + kT \ln \{ 1 - e^{-\beta \hbar \omega} \}$$

Entropy $S = \frac{U - F}{T} = k \left[\frac{\beta \hbar \omega}{e^{\beta \hbar \omega} - 1} - \ln \{ 1 - e^{-\beta \hbar \omega} \} \right]$

At absolute zero, $U \rightarrow \frac{\hbar \omega}{2}$, $C_v \rightarrow 0$ and $S \rightarrow 0$

(c) N-level System

Let the energy levels of the system be $0, \hbar\omega, 2\hbar\omega, \dots, (N-1)\hbar\omega$. Then partition function of the system is

$$\begin{aligned} Z &= \sum_r \exp(-\beta E_r) \\ &= \sum_{r=0}^{N-1} \exp(-\beta r\hbar\omega) \\ &= \frac{1 - e^{-N\beta\hbar\omega}}{1 - e^{-\beta\hbar\omega}} \end{aligned}$$

(d) Rotational Energy levels

The energies of rotational energy levels of a molecule with moment of inertia I is

$$E_J = \frac{\hbar^2}{2I} J(J+1) \quad \text{where } J = 0, 1, 2, \dots$$

are angular momentum quantum no.

Energy level E_J has degeneracy $2J+1$.

its partition function is

$$Z = \sum_{J=0}^{\infty} (2J+1) \exp(-\beta E_J)$$

$$= \sum_{J=0}^{\infty} (2J+1) \exp\left\{-\beta \hbar^2 J(J+1) / 2I\right\}$$

(e) System consists of three independent particles.

Each particle has two states of energy 0 and ϵ

Total max. energy that could be possible = 3ϵ

Min. energy that could be possible = 0

Configuration Level	Total Energy of system	No. of particles in 0 energy state	No. of particles in ϵ energy state	Degeneracy
1	0	3	0	1
2	ϵ	2	1	3
3	2ϵ	1	2	3
4	3ϵ	0	3	1

partition function of system

$$\begin{aligned} Z &= \sum_r g_r \exp(-\beta E_r) \\ &= 1 \times \exp(-\beta \cdot 0) + 3 \times \exp(-\beta \cdot \epsilon) \\ &\quad + 3 \times \exp(-\beta \cdot 2\epsilon) + 1 \times \exp(-\beta \cdot 3\epsilon) \\ &= 1 + 3e^{-\beta\epsilon} + 3e^{-2\beta\epsilon} + e^{-3\beta\epsilon} \\ &= (1 + e^{-\beta\epsilon})^3 \end{aligned}$$

The partition function for one particle system

$$Z_1 = 1 + e^{-\beta\epsilon}$$

\therefore partition function of complete system consisting of three distinguishable particles

$$Z_3 = (1 + e^{-\beta\epsilon})^3 = (Z_1)^3$$

Similarly for two distinguishable particles

$$Z_2 = (1 + e^{-\beta\epsilon})^2 = (Z_1)^2$$

probability that the system is in state of energy E_r

$$P_r = \frac{g_r \exp(-\beta E_r)}{\sum_r g_r \exp(-\beta E_r)}$$

(f) A single particle system with three energy levels of energies $0, \epsilon$ and 2ϵ . The degeneracies are $1, 3, 1$ for energy levels $0, \epsilon$ and 2ϵ respectively.

partition function of the system

$$\begin{aligned} Z_1 &= \sum_r g_r \exp(-\beta E_r) \\ &= 1 + 3e^{-\beta\epsilon} + e^{-2\beta\epsilon} \end{aligned}$$

$$\begin{aligned} \text{Mean energy of single particle } \langle E_r \rangle &= \frac{\sum_r E_r g_r \exp(-\beta E_r)}{\sum_r g_r \exp(-\beta E_r)} \\ \bar{E} &= \frac{0 + 3\epsilon e^{-\beta\epsilon} + 2\epsilon e^{-2\beta\epsilon}}{1 + 3e^{-\beta\epsilon} + e^{-2\beta\epsilon}} \end{aligned}$$

Partition function for a gas molecule (monoatomic)
(Single particle partition function)

Consider a gas molecule is confined in a space of volume V and in equilibrium at temperature T .

The hamiltonian of the molecule is

$$H(q, p) = \sum_{i=1}^3 \frac{p_i^2}{2m}$$

It is also the energy of the molecule.

The single particle partition function is

$$\begin{aligned} Z(T, V, 1) &= Z_1 = \frac{1}{h^3} \int d^3q d^3p \exp\{-\beta H(q, p)\} \\ &= \frac{V}{h^3} \iiint \exp\left\{-\frac{1}{2mkT} (p_x^2 + p_y^2 + p_z^2)\right\} dp_x dp_y dp_z \\ &= \frac{V}{h^3} (2\pi mkT)^{3/2} \end{aligned}$$

$$\therefore Z_1 = Z(T, V, 1) = \frac{V}{\lambda^3} \quad \text{where } \lambda = \frac{h}{\sqrt{2\pi m k T}}$$

is the thermal de-Broglie wavelength of the molecule. The quantity $\sqrt{2\pi m k T}$ has the character of an average thermal momentum of a molecule.

System of N non-interacting particles

Consider a system of N non-interacting particles which are confined in a space of volume V and are in equilibrium at temperature T . The hamiltonian of the system is sum of N one particle hamiltonians.

$$H(q_1, \dots, q_{3N}, p_1, \dots, p_{3N}) = \sum_{i=1}^N h(q_i, p_i)$$

$$\therefore H(q, p) = \sum_{i=1}^{3N} \frac{p_i^2}{2m} \quad \text{for } N \text{ particles.}$$

Particles of the system are distinguishable.

∴ partition function of the system

$$\begin{aligned} Z_N &= Z(T, V, N) = \frac{1}{h^{3N}} \int \cdots \int e^{-\beta \sum_{i=1}^{3N} \frac{p_i^2}{2m}} d^3p d^3q \\ &= \frac{V^N}{h^{3N}} \int \cdots \int e^{-\beta \left(\frac{p_1^2}{2m} + \frac{p_2^2}{2m} + \cdots + \frac{p_{3N}^2}{2m} \right)} dp_1 dq_1 \cdots dp_{3N} dq_{3N} \\ &= \frac{V^N}{h^{3N}} \left(\int \sqrt{2\pi m kT} \right)^{3N} \\ &= \left[\frac{V}{h^3} (2\pi m kT)^{3/2} \right]^N \\ &= [Z(T, V, 1)]^N = Z_1^N \end{aligned}$$

If the particles of the system are indistinguishable, then partition function of the system is

$$\begin{aligned} Z_N &= Z(T, V, N) = \frac{1}{N!} [Z(T, V, 1)]^N \\ &= \frac{1}{N!} \left(\frac{V}{\lambda^3} \right)^N = \frac{1}{N!} \left[\frac{V}{h^3} (2\pi m kT)^{3/2} \right]^N \end{aligned}$$

$$\text{Now, } Z(T, V, N) = \frac{1}{N!} \left[\frac{V}{h^3} (2\pi m k T)^{3/2} \right]^N$$

$$\begin{aligned} \ln Z(T, V, N) &= N \ln \left[\frac{V}{h^3} (2\pi m k T)^{3/2} \right] - \ln N! \\ &= N \ln \left[\frac{V}{h^3} (2\pi m k T)^{3/2} \right] - N \ln N + N \\ &= N \left[\ln \left[\frac{V}{N h^3} (2\pi m k T)^{3/2} \right] + 1 \right] \end{aligned}$$

\therefore Helmholtz free energy

$$\begin{aligned} F(T, V, N) &= -k T \ln \{ Z(T, V, N) \} \\ &= N k T \left[\ln \left\{ \frac{N}{V} \left(\frac{h^2}{2\pi m k T} \right)^{3/2} \right\} - 1 \right] \end{aligned}$$

Internal energy $U = - \frac{\partial \ln Z}{\partial \beta} = \frac{3}{2} N k T$

specific heat $C_V = \left(\frac{\partial U}{\partial T} \right)_V = \frac{3}{2} N k$

pressure $p = -\left(\frac{\partial F}{\partial V}\right)_{T,N} = \frac{NkT}{V}$

Entropy

$$S = -\left(\frac{\partial F}{\partial T}\right)_{N,V} = \frac{U-F}{T}$$
$$= Nk \left[\frac{5}{2} + \ln \left[\frac{V}{N} \left(\frac{2\pi m k T}{h^2} \right)^{3/2} \right] \right]$$

chemical potential

$$\mu = \left(\frac{\partial F}{\partial N}\right)_{V,T} = kT \ln \left\{ \frac{N}{V} \left(\frac{h^2}{2\pi m k T} \right)^{3/2} \right\}$$

Gibbs function

$$G = H - TS$$
$$= U + PV - TS$$
$$= F + PV$$
$$= NkT \left[\ln \left\{ \frac{N}{V} \left(\frac{h^2}{2\pi m k T} \right)^{3/2} \right\} \right]$$

References:

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Thank You

For any questions/doubts/suggestions and submission of assignments

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