Raman Spectroscopy

Part I: Introduction of Raman spectroscopy

B.Sc. (H) Chemistry

Dr Anil Kumar Singh Department of Chemistry Mahatma Gandhi Central University

Introduction

The Raman effect was named after one of its discoverers, the Indian scientist C. V. Raman, who observed the effect in organic liquids in 1928 together with K. S. Krishnan, and independently by Grigory Landsberg and Leonid Mandelstam in inorganic crystals.

- Raman won the Nobel Prize in Physics in 1930 for this discovery.
- The first observation of Raman spectra in gases was in 1929 by Franco Rasetti.



Sir Chandrasekhara Venkata Raman

Quantum Theory of Raman Effect

- Raman Scattering could be understood in terms of the quantum theory of radiation.
- > Photons can be imagined to undergo collisions with molecules.



Quantum Theory of Raman Effect

Radiation scattered with a frequency lower than that of the incident beam is referred to as Stokes' radiation, while that at higher frequency is called anti-Stokes' radiation.



Quantum Theory of Raman Effect

The molecule can gain or lose amounts of energy only in accordance with the quantal laws; i.e. its energy change, ΔE joules, must be the difference in energy between two of its allowed states.



Classical Theory of the Raman Effect: Molecular Polarizability

- The basic concept of this spectroscopy depends on the polarizability of a molecule and the applied Electric field.
- When a molecule is put into a static electric field it suffers some distortion, the positively charged nuclei being attracted towards the negative pole of the field, the electrons to the positive pole.

induced electric dipole moment

Molecule in static electric field

E

Molecule in zero electric field

Molecular Polarizability

The size of the induced dipole μ , depends both on the magnitude of the applied field, **E**, and on the ease with which the molecule can be distorted.

Polarizability is anisotropic

$$\mu = \alpha E$$

where α is the polarizability of the molecule.



• Electrons forming the bond are less easily displaced by the field across the bond axis

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Polarizability Ellipsoid

The polarizability of a molecule in various directions is conventionally represented by drawing a polarizability ellipsoid.



In general a polarizability ellipsoid is defined as a three-dimensional surface whose distance from the electrical centre of the molecule (in H₂ this is also the centre of gravity) is proportional to $1/\sqrt{\alpha}$, where α is the polarizability along the line joining a point i on the ellipsoid with the electrical centre.

Ellipsoid axis $\propto 1/\sqrt{\alpha}$

Polarizability Ellipsoid

Points to remember

- The student must not make the mistake of confusing a polarizability ellipsoid with electron orbitals or electron clouds.
- In a sense the polarizability ellipsoid is the inverse of an electron cloud-where the electron cloud is largest the electrons are further from the nucleus and so are most easily polarized.
- This, as we have seen, is represented by a small axis for the polarizability ellipsoid.

Molecular Polarizability

When a molecules is subjected to a beam of radiation of frequency ' ν ' the electric field experienced by each molecule varies according to the equation

 $E = E_0 \sin 2\pi v t$

and thus the induced dipole also undergoes oscillations of frequency v:

 $\mu = \alpha E$

or, $\mu = \alpha E_0 \sin 2\pi v t$

This is the classical explanation of **Rayleigh scattering**. Here the oscillating dipole emits radiation of its own oscillation frequency v.

If, in addition, the molecule undergoes some internal motion, such as vibration or rotation, which changes the polarizability periodically, then the oscillating dipole will have superimposed upon it the vibrational or rotational oscillation.

$$\alpha = \alpha_0 + \beta \sin 2\pi \nu_{vib} t$$
$$\mu = \alpha E$$
$$\mu = (\alpha_0 + \beta \sin 2\pi \nu_{vib} t) E_0 \sin 2\pi \nu t$$

We know the trigonometric expression,

$$\sin A \sin B = \frac{1}{2} \{ \cos(A - B) - \cos(A + B) \\ \mu = (\alpha_0 + \beta \sin 2\pi v_{vib} t) E_0 \sin 2\pi v t \\ \mu = \alpha_0 E_0 \sin 2\pi v t + \beta \sin 2\pi v_{vib} t \cdot E_0 \sin 2\pi v t \\ \mu = \alpha_0 E_0 \sin 2\pi v t + \frac{1}{2} \beta E_0 \{ \cos 2\pi (v - v_{vib}) - \cos 2\pi (v + v_{vib}) t \} \\ Rayleigh Scattering \qquad Anti stokes line \qquad stokes line \qquad stokes line \qquad description and the stokes line \qquad description a$$

Molecular Polarizability

In order to be Raman active a molecular rotation or vibration must cause some change in a component of the molecular polarizability. A change in polarizability is, of course, reflected by a change in either the *magnitude* or the *direction* of the polarizability ellipsoid.

Reference Fundamentals of Molecular Spectroscopy, 4th Ed. By Colin N. Banwell and Elaine M. McCash

Next: Rotational & Vibrational Raman Spectra

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